

Własności oktupolowej przestrzeni kolektywnej

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- ▶ Wprowadzenie
- ▶ Własności przestrzeni oktupolowej
 - ▶ Układ wewnętrzny
 - ▶ Operatory: moment pędu, hamiltonian
- ▶ Uwagi końcowe

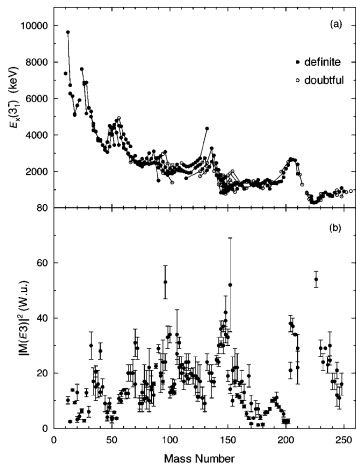
- ▶ Aparatura
- ▶ Oktupole w Warszawie i okolicach. Długa historia
 - ▶ J. Chwaszczewska, R. Kaczarowski, J. Rudzińska, W. Kurcewicz, and J. Żylicz. *Possible $K^\pi = 2^-$ and $K^\pi = 1^-$ Octupole Bands in ^{232}U* . Contrib. Intern. Symp. Nucl. Struct., Dubna, p.59 (1968).
 - ▶ J. Błocki and W. Kurcewicz. *Octupole Vibrations of Even Nuclei in the Transuranic Region*. Phys. Lett. B **30** (1969), 458.
 - ...
 - ▶ L. P. Gaffney et al., Nature, **497** (2013) 199
 - ...
- ▶ Doświadczenie
- ▶ Teoria
 1. W stronę fizyki
 2. W stronę matematyki

2. W stronę matematyki

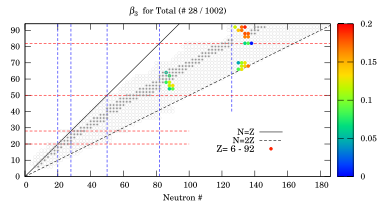
S.G. Rohoziński, *Oscillator basis for octupole collective motion in nuclei*,
J.Phys. G Nuclear and Particle Physics, **4** (1978) 1075

...

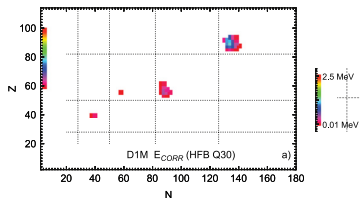
S.G. Rohoziński and L.P., *The octupole collective Hamiltonian. Does it follow the example of the quadrupole case?*, Walter Greiner Memorial Volume (2018).



T. Kibedi & H. Spears, *At.Data Nucl.Data* **80**, 35–82 (2002)



S. Ebata and T. Nakatsukasa, 2017, *Phys. Scr.* **92** 064005



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7-wymiarowa nieprzywiedlna reprezentacja grupy $O(3)$, $P = -I$
Zmienne zespolone $\alpha_{3\mu} = \alpha_\mu, \mu = -3, \dots, 3$ z warunkiem

$$\alpha_{-\mu} = (-)^{\mu} \alpha_{\mu}^*$$

Działanie grupy $SO(3)$, $\omega = (\phi, \theta, \psi)$ — kąty Eulera

$$\tilde{\alpha}_{\mu} = \sum_{\nu} \mathcal{D}_{\mu\nu}^3(\omega) \alpha_{\nu}, \quad \mathcal{D}_{\mu\nu}^3(\omega) = e^{i\mu\phi} d_{\mu\nu}^3(\theta) e^{i\nu\psi}$$

Dwa podstawowe przykłady

1.

$$r(\theta, \phi) = r_0 \left(1 + \sum_{\mu} \alpha_{\mu} Y_{3\mu}^*(\theta, \phi) \right)$$

2.

$$\alpha_{\mu} \sim \langle Q_{3\mu} \rangle, \quad Q_{3\mu} \sim \sum_i r_i^3 Y_{3\mu}(\theta_i, \phi_i)$$

Zmienne rzeczywiste $a_0, a_m, b_m, m = 1, 2, 3$

$$\alpha_\mu = (a_\mu + ib_\mu)/\sqrt{2}$$

$$\alpha_0 = a_0$$

$$\alpha_{-\mu} = (-)^\mu (a_\mu - ib_\mu)/\sqrt{2}$$

Działanie grupy SO(3)

$$\tilde{a}_m = \sum_{k<0} T_{mk}^{++}(\boldsymbol{\omega}) a_k + \sum_{k>0} T_{mk}^{-+}(\boldsymbol{\omega}) b_k$$

$$\tilde{b}_m = \sum_{k<0} T_{mk}^{+-}(\boldsymbol{\omega}) a_k + \sum_{k>0} T_{mk}^{--}(\boldsymbol{\omega}) b_k$$

$$T_{mk}^{\pm\pm}(\boldsymbol{\omega}) \sim \begin{Bmatrix} \cos m\phi \\ \sin m\phi \end{Bmatrix} d_{mk}^3(\theta) \begin{Bmatrix} \cos m\psi \\ \sin m\psi \end{Bmatrix}$$

Jeszcze raz działanie grupy $SO(3)$

$$\tilde{\alpha}_\mu = D_{\mu 0}^{(+)}(\boldsymbol{\omega}) a_0 + \sum_{k>0} \left(D_{\mu k}^{(+)}(\boldsymbol{\omega}) a_k + D_{\mu k}^{(-)}(\boldsymbol{\omega}) b_k \right)$$

$$D_{\mu 0}^{(+)} = \mathcal{D}_{\mu 0}^3$$

$$D_{\mu k}^{(+)} = (\mathcal{D}_{\mu k}^3 + (-)^k \mathcal{D}_{\mu, -k}^3) / \sqrt{2}$$

$$D_{\mu k}^{(-)} = i(\mathcal{D}_{\mu k}^3 - (-)^k \mathcal{D}_{\mu, -k}^3) / \sqrt{2}$$

$D_{\mu k}^{(+)}, D_{\mu k}^{(-)}$ tworzą macierz unitarną (tak jak $\mathcal{D}_{\mu\mu'}$)

Analogiczne definicje dla kwadrupoli i innych multipoli

- ▶ Elipsoida, $\sum_{i,j=1}^3 B_{ij}x_i x_j = 0$
- ▶ Tensory jako symetryczne, bezśladowe macierze kwadratowe 3×3 , $\tilde{W} = RWR^T$
- ▶ Układ wewnętrzny: układ osi głównych elipsoidy, postać diagonalna macierzy symetrycznej

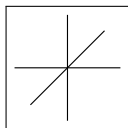
$$(\tilde{a}_{2m}, \tilde{b}_{2m}) \rightarrow (a_{20}, a_{22}, \omega)$$

$$\omega \leftarrow a_{21} = b_{21} = b_{22} = 0 \text{ czyli } \alpha_{\pm 1} = 0, \alpha_2 = \alpha_{-2}$$

$$a_{20} = \beta \cos \gamma, \quad a_{22} = \beta \sin \gamma$$

Niezmienność ze względu na grupę O_h

Nieprzywiedlne reprezentacje O_h



IR	A_1^\pm	A_2^\pm	E^\pm	F_1^\pm	F_2^\pm
dim	1	1	2	3	3

QC

$$E^+ : a_{20}, a_{22}$$

$$F_2^+ : a_{21}, b_{21}, b_{22}$$

Przestrzeń oktopolowa

$$A_2^- : b = b_3$$

$$F_1^- : f_x = \sqrt{3/8}a_1 - \sqrt{5/8}a_3, \quad f_y = \sqrt{3/8}b_1 + \sqrt{5/8}b_3, \quad f_z = a_0$$

$$F_2^- : g_x = \sqrt{5/8}a_1 + \sqrt{3/8}a_3, \quad g_y = -\sqrt{5/8}b_1 + \sqrt{3/8}b_3, \quad g_z = a_2$$

$$\alpha_\mu \rightarrow (a_m, b_m) \rightarrow (b, f_s, g_s)$$

$$\tilde{\alpha}_\mu = A_\mu(\boldsymbol{\omega})b + \sum_{s=x,y,z} [F_{\mu s}(\boldsymbol{\omega})f_s + G_{\mu s}(\boldsymbol{\omega})g_s]$$

Funkcje „kubiczne”

$$A_\mu(\boldsymbol{\omega}) = D_{\mu 2}^{(-)}(\boldsymbol{\omega})$$

$$F_{\mu x}(\boldsymbol{\omega}) = \sqrt{3/8} D_{\mu 1}^{(+)}(\boldsymbol{\omega}) - \sqrt{5/8} D_{\mu 3}^{(+)}(\boldsymbol{\omega}),$$

$$F_{\mu y}(\boldsymbol{\omega}) = \sqrt{3/8} D_{\mu 1}^{(-)}(\boldsymbol{\omega}) + \sqrt{5/8} D_{\mu 3}^{(-)}(\boldsymbol{\omega}),$$

$$F_{\mu z}(\boldsymbol{\omega}) = D_{\mu 0}^{(+)}(\boldsymbol{\omega})$$

$$G_{\mu x}(\boldsymbol{\omega}) = \sqrt{5/8} D_{\mu 1}^{(+)}(\boldsymbol{\omega}) + \sqrt{3/8} D_{\mu 3}^{(+)}(\boldsymbol{\omega})$$

$$G_{\mu y}(\boldsymbol{\omega}) = -\sqrt{5/8} D_{\mu 1}^{(-)}(\boldsymbol{\omega}) + \sqrt{3/8} D_{\mu 3}^{(-)}(\boldsymbol{\omega})$$

$$G_{\mu z}(\boldsymbol{\omega}) = D_{\mu 2}^{(+)}(\boldsymbol{\omega}).$$

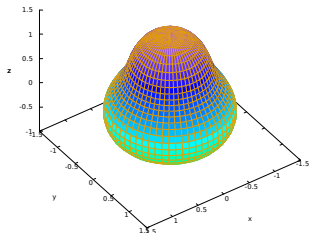
Oktupoloida

$$r(\theta, \phi) = r_0 \left(1 + \sum_{\mu} \alpha_{\mu} Y_{3\mu}^*(\theta, \phi) \right) = r_0 \left(1 + b v_b(\theta, \phi) + \sum_s f_s v_{f_s}(\theta, \phi) + \sum_s g_s v_{g_s}(\theta, \phi) \right)$$

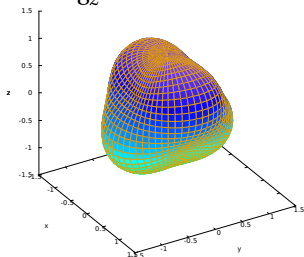
Przykład

$$v_{f_x} = \sqrt{2} \left(\sqrt{3/8} \operatorname{Re} Y_{31} - \sqrt{5/8} \operatorname{Re} Y_{33} \right) = \frac{\sqrt{7} x (3y^2 + 3z^2 - 2x^2)}{4\sqrt{\pi} r^3} \sim F_{0x}$$

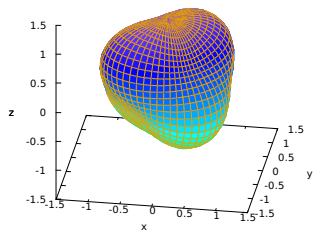
$f_z = 0.4$



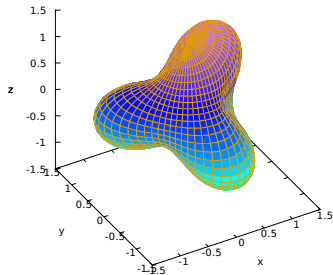
$g_z = 0.4$



$$b = 0.4$$



$$b = f_s = g_s = 0.3$$



$$A_2^- \oplus F_1^- \oplus F_2^- \\ \alpha_\mu \rightarrow (a_m, b_m) \rightarrow (b, f_s, g_s)$$

Układ wewnętrzny dla oktapoli, dwa warianty:

$$\alpha_\mu \rightarrow \begin{cases} \mathcal{F}: & (b, f_s, \omega), & \omega \leftarrow g_s = 0 \\ \mathcal{G}: & (b, g_s, \omega), & \omega \leftarrow f_s = 0 \end{cases}$$

Dla \mathcal{F}

$$\alpha_\mu = A_\mu(\omega)b + \sum_{s=x,y,z} F_{\mu s}(\omega)f_s$$

Czy warianty \mathcal{F} , \mathcal{G} są równoważne?

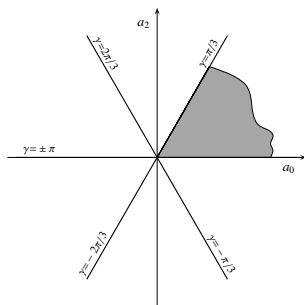
„Przejście do układu wewnętrznego" — zamiana zmiennych

Wybrane ogólne własności zmiennych wewnętrznych

- ▶ Moment pędu — funkcja ω, ∂_ω
- ▶ Energia potencjalna nie zależy od ω

Jakobian, jednoznaczność, minimalny zakres

$$|\det Z_{QC}| = 2|(a_{22}^2 - 3a_{20}^2)a_{22}| \sin\theta = 8\beta^3 \sin 3\gamma \sin\theta$$

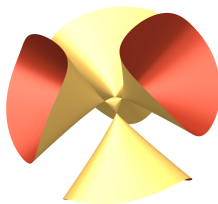
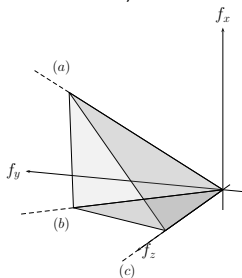


Jakobian

$$|\det J_{\mathcal{F}}| = 8 \left(b^3 - 15b(f_x^2 + f_y^2 + f_z^2)/16 + 15\sqrt{15}f_x f_y f_z/32 \right) \sin \theta = d_f \sin \theta$$

$$|\det J_{\mathcal{G}}| = 15\sqrt{15}g_x g_y g_z \sin \theta / 4 = d_g \sin \theta$$

Jednoznaczność, minimalny zakres zmiennych b, f_s, g_s



$$\det J_{\mathcal{F}} = 0, b = 0.5$$

Ogólnie $[\partial_{q_i}], \quad \tilde{q}_j = \tilde{q}_j(\mathbf{q}) \quad \partial_{\tilde{q}_j} = \sum_k \frac{\partial q_k}{\partial \tilde{q}_j} \partial_{q_k}$

$$\sum c_{ij}(\mathbf{q}) \partial_{q_j} \stackrel{?}{=} \partial_{\tilde{q}_i}$$

Oktupole $[\partial_b, \partial_s, \partial_\omega] \rightarrow [\partial_b, \partial_s, S_k], \quad \partial_s = \partial_{f_s} \text{ lub } \partial_{g_s}$

$$S_x = -\frac{\cos \psi}{\sin \theta} \partial_\phi + \sin \psi \partial_\theta + \cos \psi \operatorname{ctg} \theta \partial_\psi$$

$$S_y = \frac{\sin \psi}{\sin \theta} \partial_\phi + \cos \psi \partial_\theta - \sin \psi \operatorname{ctg} \theta \partial_\psi$$

$$S_z = \partial_\psi$$

S_k — generatory prawej regularnej reprezentacji grupy $SO(3)$, $(T_R f)(S) = f(SR)$

Operatory fizyczne $S_k^{(\text{ph})} = iS_k$

Zmienne (a_m, b_m)

$$g_R = I_7 \quad \rightarrow \quad \sum_m \dot{a}_m^2 + \sum_m \dot{b}_m^2$$

Zmienne α_μ

$$(g_C)_{\mu\nu} = (-1)^\mu \delta_{\mu,-\nu} \quad \rightarrow \quad \sum_\mu \dot{\alpha}_\mu \dot{\alpha}_\mu^*$$

Przypadek \mathcal{F} , $(\dot{b}, \dot{f}_s, \boldsymbol{\eta})$

$$g_{\mathcal{F}} = \begin{pmatrix} I_4 & M_{12} \\ M_{12}^T & M_{22} \end{pmatrix}$$

$$M_{12} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -3f_z/2 & 3f_y/2 \\ 3f_z/2 & 0 & 3f_x/2 \\ -3f_y/2 & 3f_x/2 & 0 \end{pmatrix}$$

Kwadrupole $(\dot{a}_{20}, \dot{a}_{22}, \boldsymbol{\eta})$

$$g_{\text{quad}} = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & (a_{22} + \sqrt{3}a_{20})^2 & 0 & 0 \\ 0 & 0 & 0 & (a_{22} - \sqrt{3}a_{20})^2 & 0 \\ 0 & 0 & 0 & 0 & 4a_{22}^2 \end{pmatrix}$$

Nowa baza w przestrzeni stycznej

$$W_{k,f} = S_k + \Lambda_{k,f}$$

$$\Lambda_{1,f} = 3(f_y \partial_{f_z} - f_z \partial_{f_y})/2$$

$$\Lambda_{2,f} = 3(f_z \partial_{f_x} - f_x \partial_{f_z})/2$$

$$\Lambda_{3,f} = 3(f_x \partial_{f_y} - f_y \partial_{f_x})/2$$

Teraz

$$g_{\mathcal{F}} = \begin{pmatrix} I_4 & M_{12} \\ M_{12}^T & M_{22} \end{pmatrix} \rightarrow g'_{\mathcal{F}} = \begin{pmatrix} I_4 & 0 \\ 0 & T_f \end{pmatrix}$$

$$T_f = \frac{1}{4} \begin{pmatrix} 16b^2 + 15f_y^2 + 15f_z^2 & 15f_x f_y + 8\sqrt{15}bf_z & 8\sqrt{15}bf_y + 15f_x f_z \\ 15f_x f_y + 8\sqrt{15}bf_z & 16b^2 + 15f_x^2 + 15f_z^2 & 8\sqrt{15}bf_x + 15f_y f_z \\ 8\sqrt{15}bf_y + 15f_x f_z & 8\sqrt{15}bf_x + 15f_y f_z & 16b^2 + 15f_x^2 + 15f_y^2 \end{pmatrix}$$

Przypadek \mathcal{F}

$$\alpha_\mu = A_\mu(\boldsymbol{\omega})b + \sum_{s=x,y,z} F_{\mu s}(\boldsymbol{\omega})f_s$$

$$\sum_\mu G_{\mu s}^*(\boldsymbol{\omega})\alpha_\mu = 0 \quad (g_s)$$

$$b = \sum_\mu A_\mu^*(\boldsymbol{\omega})\alpha_\mu$$

$$f_s = \sum_\mu F_{\mu s}^*(\boldsymbol{\omega})\alpha_\mu$$

$$\frac{\partial}{\partial \alpha_\mu} = A_\mu^* \frac{\partial}{\partial b} + \sum_s F_{\mu s}^* \frac{\partial}{\partial f_s} + \frac{1}{d_f} \sum_k x_{k,f}^* W_{k,f}$$

$$x_{1,f} = G_{\mu x} \left(\frac{15f_x^2}{4} - 4b^2 \right) + G_{\mu y} \left(\sqrt{15}bf_z - \frac{f_x f_y}{4} \right) + G_{\mu z} \left(\sqrt{15}bf_y - \frac{f_x f_z}{4} \right)$$

...

Operator Laplace'a (-Beltramiego)

W bazie współrzędnościowej

$$\Delta W = \frac{1}{\sqrt{\det g}} \sum_{ij} \frac{\partial}{\partial q_i} \sqrt{\det g} g^{ij} \frac{\partial W}{\partial q_j}$$
$$\sum_{\mu} \frac{\partial}{\partial \alpha_{\mu}} \frac{\partial}{\partial \alpha_{\mu}^*} = \frac{\partial^2}{\partial a_0^2} + \sum_{m=1,2,3} \left(\frac{\partial^2}{\partial a_m^2} + \frac{\partial^2}{\partial b_m^2} \right)$$

$$\Delta f = \operatorname{div} \operatorname{grad} f$$

$$\operatorname{div} Y = \sum_k (X_k Y^k + \sum_j \Gamma_{kj}^j Y^k) = \sum_k (X_k + \Omega_k) Y^k$$

X_k — wektory bazowe w przestrzeni stycznej, $\Omega_k = \sum_j \Gamma_{kj}^j$

Dla baz współrzędnościowych $X_k = \partial_{q_k}$ i $\Omega_k = \partial_{q_k} \ln \sqrt{\det g}$

Ogólniej

$$\Gamma_{ijk} = \frac{1}{2} (g_{ij,k} + g_{ik,j} - g_{jk,i} + c_{ijk} + c_{ikj} - c_{jki})$$

$$g_{ij,k} = X_k g_{ij}$$

$$[X_i, X_j] = \sum_m c_{ij}^m X_m, \quad c_{ijk} = \sum_m g_{km} c_{ij}^m$$

$$\Gamma_{kj}^m = (g^{-1})^{mi} \Gamma_{ikj}$$

$$\sum_{\mu} \frac{\partial}{\partial \alpha_{\mu}} \frac{\partial}{\partial \alpha_{\mu}^*} = \frac{1}{d_f} \frac{\partial}{\partial b} d_f \frac{\partial}{\partial b} + \sum_s \frac{1}{d_f} \frac{\partial}{\partial f_s} d_f \frac{\partial}{\partial f_s} + \frac{1}{d_f} \sum_{k,j} W_{k,f} M_{kj}^{(f)} d_f W_{j,f}$$

$$d_f = 8 \left(b^3 - 15b(f_x^2 + f_y^2 + f_z^2) / 16 + 15\sqrt{15}f_x f_y f_z / 32 \right)$$

$$W_{k,f} = S_k + \Lambda_{k,f}$$

$$\Lambda_{1,f} = 3(f_y \partial_{f_z} - f_z \partial_{f_y}) / 2$$

...

$$M^{(f)} = T_f^{-1}$$

$$L_{1\kappa}(\boldsymbol{\alpha}, \boldsymbol{\partial}\boldsymbol{\alpha}) = -2\sqrt{7} \sum_{\mu\nu} (3\mu 3\nu | 1\kappa) \alpha_\mu \frac{\partial}{\partial \alpha_\nu^*}$$

Zmienne wewnętrzne

$$L_{1\kappa} = \mathcal{D}_{\kappa 0}^1(\boldsymbol{\omega}) S_z^{(\text{ph})} + \mathcal{D}_{\kappa 1}^1(\boldsymbol{\omega}) \left(-S_x^{(\text{ph})} + iS_y^{(\text{ph})} \right) / \sqrt{2} + \mathcal{D}_{\kappa, -1}^1(\boldsymbol{\omega}) \left(S_x^{(\text{ph})} + iS_y^{(\text{ph})} \right) / \sqrt{2}$$

$S_{x,y,z}^{(\text{ph})}$ zależą tylko od $(\boldsymbol{\omega}, \boldsymbol{\partial}\boldsymbol{\omega})$

To nie przypadek

$$\alpha_\mu = \sum_\nu \mathcal{D}_{\mu\nu}(\Omega^{-1}) \tilde{\alpha}_\nu$$

$$b = \sum_\mu A_\mu^*(\boldsymbol{\omega}) \alpha_\mu = \sum_{\mu,\nu} A_\mu^*(\boldsymbol{\omega}) \mathcal{D}_{\mu\nu}(\Omega^{-1}) \tilde{\alpha}_\nu = \sum_\nu A_\nu^*(\boldsymbol{\omega} \circ \Omega) \tilde{\alpha}_\nu$$

$$f_s = \sum_\mu F_{\mu s}^*(\boldsymbol{\omega}) \alpha_\mu = \dots = \sum_\nu F_{\nu s}^*(\boldsymbol{\omega} \circ \Omega) \tilde{\alpha}_\nu$$

$$0 = \sum_\mu G_{\mu s}^*(\boldsymbol{\omega}) \alpha_\mu = \dots = \sum_\nu G_{\nu s}^*(\boldsymbol{\omega} \circ \Omega) \tilde{\alpha}_\nu \quad \longrightarrow \quad (b, f_s, \boldsymbol{\omega}) \rightarrow (b, f_s, \tilde{\boldsymbol{\omega}} = \boldsymbol{\omega} \circ \Omega)$$

Elementarne tensory

$$\mathbf{t}_l^{(n)} = \underbrace{[\boldsymbol{\alpha} \times \dots \times \boldsymbol{\alpha}]_l}_n$$

QC. Elementarne skalary $l=0, n=2,3$

$$[\boldsymbol{\alpha}_2 \times \boldsymbol{\alpha}_2]_0 \sim \beta^2, \quad [[\boldsymbol{\alpha}_2 \times \boldsymbol{\alpha}_2]_2 \times \boldsymbol{\alpha}_2]_0 \sim \beta^3 \cos 3\gamma$$

Oktupole. Elementarne skalary, $l=0, n=2,4,6,10,15$

$$\mathbf{t}_0^{(2)} = -\frac{1}{\sqrt{7}}(b^2 + \sigma_2)$$

$$\mathbf{t}_0^{(4)} = \frac{1}{84\sqrt{5}}(16\sigma_4 - 13\sigma_{42} + 80b^2\sigma_2 + 24\sqrt{15}b\sigma_3)$$

$$\sigma_2 = f_x^2 + f_y^2 + f_z^2, \quad \sigma_4 = f_x^4 + f_y^4 + f_z^4, \quad \sigma_6 = f_x^6 + f_y^6 + f_z^6$$

$$\sigma_3 = f_x f_y f_z, \quad \sigma_{42} = f_x^2 f_y^2 + f_y^2 f_z^2 + f_x^2 f_z^2$$

$$H(\boldsymbol{\alpha}, \boldsymbol{\partial\alpha}) = -\frac{1}{2W(\boldsymbol{\alpha})} \sum_{\mu, \nu} \frac{\partial}{\partial\alpha_\mu} W(\boldsymbol{\alpha}) B_{\mu\nu}^{-1}(\boldsymbol{\alpha}) \frac{\partial}{\partial\alpha_\nu} + V(\boldsymbol{\alpha}),$$

zwykle $W(\boldsymbol{\alpha}) = \sqrt{\det(B_{\mu\nu}(\boldsymbol{\alpha}))}$.

$B_{\mu\nu}(\boldsymbol{\alpha})$

- ▶ symetryczny i dodatnio określony (w zmiennych rzeczywistych)
- ▶ można go rozłożyć na tensory

$$B_{\mu\nu}(\boldsymbol{\alpha}) = \sum_{\lambda=0,2,4,6} (3\mu 3\nu | \lambda \kappa) \tau_{\lambda\kappa}(\boldsymbol{\alpha})$$

- ▶ 28 niezależnych składowych (nawet w zmiennych wewnętrznych)

- ▶ Kwadrupol plus oktupol
- ▶ Proste, nietrywialne przypadki
- ▶ Obliczenia mikroskopowe

- ▶ Nature, 497 (2013) 199, *Studies of pear-shaped nuclei using accelerated radioactive beams*

Going pear-shaped

